

$+1 +j/2$	$-1 -j/2$	$+3/2 +j/2$	$-1/2 +3/2$	$1/2 -1/2$	$3/10$	0	0	0	0
$+2-1/2$	$1/7$	$16/35$	$2/5$	$7/2$	$5/2$	$3/2$	$1/2$	$+3/2 -3/2$	$1/20$
$+1+1/2$	$4/7$	$1/35-2/5$	$+1/2 +1/2$	$+1/2$	$+1/2$	$+1/2$	$+1/2$	$+1/2 -1/2$	$1/4 -1/20 -1/4$
$0+3/2$	$2/7-18/35$	$1/5$	$+1/2 +1/2$	$+1/2$	$+1/2$	$+1/2$	$+1/2$	$-1/2 +1/2$	$9/20 -1/4 -1/20 1/4$
2×2	4	4	3	2	1	0	0	$-3/2 +3/2$	$1/20 -1/4 -9/20 1/4$
$+2+2$	0	$+3$	$+3$	$+2$	$+1$	0	0	$+1/2 -1/2$	$1/4 -1/20 -1/4$
$+1+1$	$1/2$	$1/2$	$1/2$	$1/2$	$1/2$	$1/2$	$1/2$	$-1/2 -1/2$	$9/20 -1/4 -1/20 1/4$
$+1+1$	0	$+2$	$3/14$	$-1/2$	$1/2$	$1/2$	$1/2$	$-1/2 +1/2$	$9/20 -1/4 -1/20 1/4$
$+1+1$	$4/7$	0	$-3/7$	4	3	2	1	$+1/2 -1/2$	$1/4 -1/20 -1/4$
$+1+1$	$3/14$	$-1/2$	$2/7$	0	0	0	0	$-1/2 +3/2$	$1/20 -1/4 -9/20 1/4$
$+2-1$	$1/14$	$3/10$	$3/7$	$1/5$	0	0	0	$+1/2 -3/2$	$2/5 1/10$
$+1+0$	$3/7$	$1/5$	$-1/5$	$-1/14$	$3/10$	0	0	$-1/2 -1/2$	$9/20 -1/4 -1/20 1/4$
$0+1$	$3/7$	$-1/5$	$-1/5$	$-1/14$	$3/10$	0	0	$-1/2 +1/2$	$1/4 -1/20 -1/4$
$-1+2$	$1/14-3/10$	$3/7$	$-1/5$	$-1/14$	$3/10$	0	0	$-2 +3/2$	$1/35-6/35$
$d_{3/2,3/2}^{3/2}$	$\frac{1+\cos\theta}{2}\cos\frac{\theta}{2}$	$d_{3/2,1/2}^2 = -\sqrt{3}\frac{1+\cos\theta}{2}\sin\frac{\theta}{2}$	$d_{2,2}^2 = \left(\frac{1+\cos\theta}{2}\right)^2$	$d_{3/2,-1/2}^{3/2} = \sqrt{3}\frac{1-\cos\theta}{2}\cos\frac{\theta}{2}$	$d_{3/2,-1/2}^2 = -\frac{1+\cos\theta}{2}\sin\theta$	$d_{3/2,-3/2}^{3/2} = -\frac{1-\cos\theta}{2}\sin\theta$	$d_{3/2,-3/2}^2 = \frac{\sqrt{6}}{4}\sin^2\theta$	$d_{1,1}^2 = \frac{1+\cos\theta}{2}(2\cos\theta-1)$	$d_{1,1}^2 = \frac{1+\cos\theta}{2}(2\cos\theta-1)$
$d_{1/2,1/2}^{3/2}$	$\frac{3\cos\theta-1}{2}\cos\frac{\theta}{2}$	$d_{1/2,-1}^2 = -\frac{1-\cos\theta}{2}\sin\theta$	$d_{1/2,-1}^2 = -\frac{1-\cos\theta}{2}\sin\theta$	$d_{1/2,0}^2 = -\sqrt{\frac{3}{2}}\sin\theta\cos\theta$	$d_{1/2,0}^2 = -\sqrt{\frac{3}{2}}\sin\theta\cos\theta$	$d_{1/2,-1}^2 = -\frac{1-\cos\theta}{2}\sin\theta$	$d_{1/2,-1}^2 = -\frac{1-\cos\theta}{2}\sin\theta$	$d_{1/2,0}^2 = \frac{1+\cos\theta}{2}(2\cos\theta-1)$	$d_{1/2,0}^2 = \frac{1+\cos\theta}{2}(2\cos\theta-1)$

$$L_- = L_{1-} + L_{2-}$$

$$\sqrt{|43\rangle} = \frac{1}{\sqrt{2}} [|2(1)122\rangle + |22\rangle |21\rangle]$$

$$\sqrt{|33\rangle} = \cancel{a |21\rangle |22\rangle} + \cancel{b |22\rangle |21\rangle}$$

$$\sqrt{a^2+b^2}=1 \quad \boxed{a=-b} \quad \boxed{\sqrt{2}}$$

$$\langle \underline{|21\rangle |22\rangle} | \underline{|21\rangle |22\rangle} \rangle = 1$$

$$|\langle \underline{|21\rangle |22\rangle} | \underline{|22\rangle |21\rangle} \rangle\rangle = 0$$

Exchange forces

Not really forces, consequence
of sym of states

1D $\psi_a(x), \psi_b(x)$, etc ^{single} _{states}

$$\frac{\psi(x_1, x_2) = \psi_a(x_1)\psi_b(x_2)}{1 \leftrightarrow \psi_b(x_1)\psi_a(x_2)}$$

$$\checkmark \psi_D(x_1, x_2) = \psi_a(x_1)\psi_b(x_2) \quad \checkmark$$

$$\checkmark \psi_+(x_1, x_2) = \frac{1}{\sqrt{2}} [\psi_a(x_1)\psi_b(x_2) + \psi_b(x_1)\psi_a(x_2)]$$

$$\checkmark \psi_-(x_1, x_2) = \frac{1}{\sqrt{2}} [\psi_a(x_1)\psi_b(x_2) - \psi_b(x_1)\psi_a(x_2)]$$

Expectation value of the distance squared btw the particles

$$\langle (x_1 - x_2)^2 \rangle = \langle x_1^2 \rangle + \langle x_2^2 \rangle - 2 \langle x_1 x_2 \rangle$$

Algebra --

$$\langle (x_1 - x_2)^2 \rangle_+ = \langle (x_1 - x_2)^2 \rangle_{\text{distinguishable}}$$
$$= 2 \left| \int_x \psi_a^*(x) \psi_b(x) dx \right|^2$$

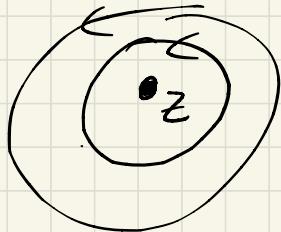
positive #

Symm state Smaller than

$$\langle (x_1 - x_2)^2 \rangle_+ < \langle (x_1 - x_2)^2 \rangle_-$$

Atoms

He $2p\ 2n$



$$Z=2$$

3-body
Nucleus is stationary approximation

Two electrons H

$$H = -\frac{\hbar^2}{2m} \left(\nabla_1^2 + \nabla_2^2 \right) - \frac{e^2}{4\pi\epsilon_0} \left[\frac{2}{r_1} + \frac{2}{r_2} \right]$$

~~$$+ \frac{e^2}{4\pi\epsilon_0} \frac{1}{|\vec{r}_1 - \vec{r}_2|}$$~~

A crossed-out term from the full Hamiltonian, showing the electron-electron interaction potential.

Ignoring interaction

$$\psi(\vec{r}_1 \vec{r}_2) = \psi_{nem}(\vec{r}_1) \psi_{n'e'm'}(\vec{r}_2)$$

ψ_{nem} are H-wavefunctions but with

$Z=2$ instead of $Z=1$

"Bohr radius" is $\frac{1}{2}$ of H-atom a_0
Energies proportional to Z^2

$$E = 4(E_n + E_{n'}) \quad E_n = -\frac{R}{n^2}$$

$$R = R_{\text{Bohr}} = 13.6 \text{ eV} \quad (\text{with highlighter})$$

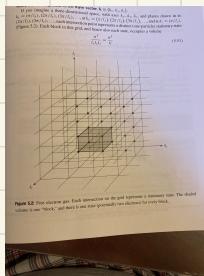
Ground state

$$\Psi_0 = \Psi_{100}(\vec{r}_1) \Psi_{100}(\vec{r}_2) = \frac{8}{\pi a_0^3} e^{-2(r_1+r_2)/a_0}$$

$$E = 4 \left(-\frac{13.6}{1} - \frac{13.6}{1} \right) = -\underline{\underline{109 \text{ eV}}}$$

The value is measured $E = -79 \text{ eV}$

30% correction from interaction
between the two electrons



H states

~~uuuuu~~ $n=3 \quad l=2, 1, 0 \quad (9) \quad 10$

~~ooooo~~ $n=2 \quad l=1, 0 \quad m=0, \pm 1, m=0 \quad (4)$

~~o~~ $n=1 \quad l=0 = m=0 \quad (1)$

n^2 = degeneracy

electrons also have spin degree
of freedom $\Rightarrow \times 2$

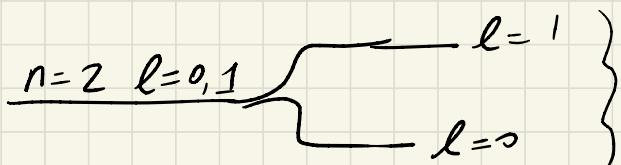
Contact with Chemistry

n = shell Chemistry

Naive picture only works for $n=1, 2$

Occupancy in $n^2 = 1 \times 2 = 2$
 $n^2 = 2^2 = 4 \times 2 = 8$
 $n^2 = 3^2 = 9 \times 2 = 18$

Naive
picture



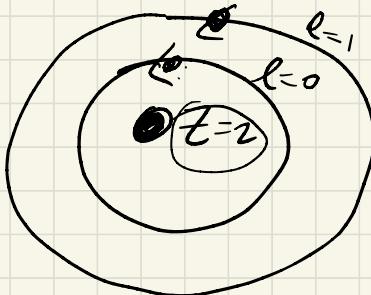
interaction between electron with some n breaks the degeneracy

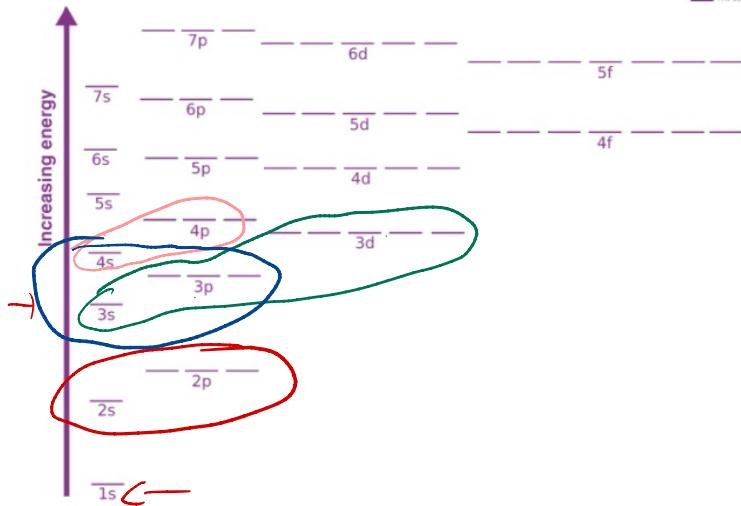
$$\textcircled{a} \quad E(n) \rightarrow E(n, l)$$

Hand writing:

$l=0$ typical radius is smaller

than in $l=1$





1S $n=1 \quad l=0$

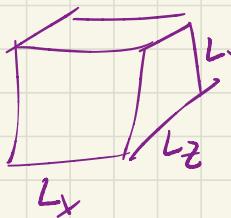
$l=0$ S
 $l=1$ P
 $l=2$ D
 $l=3$ F

SOLIDS

Electrons on outer shells loosely bound - simplest model is that there are some electrons that are free

Treat electrons as free from being bound in the atom, and not interacting with each other

Electron gas (free)



$$V(x) = 0 \text{ inside}$$
$$0 < x < L_x$$
$$0 < y < L_y$$
$$0 < z < L_z$$

$$V(x) = \infty \text{ outside}$$

Particle in a box

$$\psi_{n_1, n_2, n_3} = \sqrt{\frac{8}{L_x L_y L_z}} \sin \frac{n_1 \pi x}{L_x} \sin \frac{n_2 \pi y}{L_y} \sin \frac{n_3 \pi z}{L_z}$$

$$E = \frac{\pi^2 k^2}{m}$$

$$\vec{k} = (k_x, k_y, k_z)$$

$$k_x = \frac{n_x \pi}{L_x} \quad k_y = \frac{n_y \pi}{L_y} \quad k_z = \frac{n_z \pi}{L_z}$$

If there was no box, free particles

Wavefunction for free particle is just

a wave $e^{i \frac{\vec{p} \cdot \vec{r}}{\hbar}}$ \vec{p} is the momentum

For free particle \vec{p} takes any value

$\vec{k} = \frac{\vec{p}}{\hbar}$ takes any value

In the box you have (almost) free particle states but with only some values of k that are allowed

If electrons were bosons what

would be the ground state energy of the system?

$E = \frac{\pi^2 k^2}{m}$ lowest energy for one electron is

$$k = \frac{\pi}{L_x}$$

All electrons will pile up

In the lowest state

Electrons are fermions

⇒ They cannot all pile up in the lowest energy state

where \vec{k} is the wave vector, $\vec{k} \equiv (k_x, k_y, k_z)$.

If you imagine a three-dimensional space, with axes k_x , k_y , k_z and planes drawn in at $k_x = (\pi/l_x), (2\pi/l_x), (3\pi/l_x), \dots$, at $k_y = (\pi/l_y), (2\pi/l_y), (3\pi/l_y), \dots$, and at $k_z = (\pi/l_z), (2\pi/l_z), (3\pi/l_z), \dots$, each intersection point represents a distinct (one-particle) stationary state (Figure 5.2). Each block in this grid, and hence also each state, occupies a volume

$$\frac{\pi^3}{l_x l_y l_z} = \frac{\pi^3}{V} \quad (5.51)$$

6

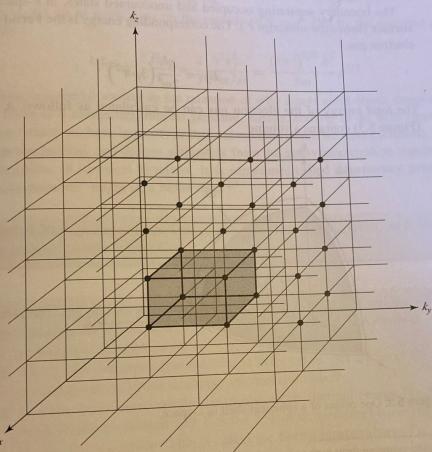


Figure 5.2: Free electron gas. Each intersection on the grid represents a stationary state. The shaded volume is one "block," and there is one state (potentially two electrons) for every block.