

$$\vec{J} = -e N_e \vec{V}$$

↑ ↓

must increase as one goes across
decrease to compensate, when $\frac{\partial V}{\partial t} = 0$.

Ohm's Law

Empirically, in many substances,

$$\vec{J} \propto \vec{E}, \quad \vec{J} = \sigma \vec{E} = \frac{1}{\rho} \vec{E}$$

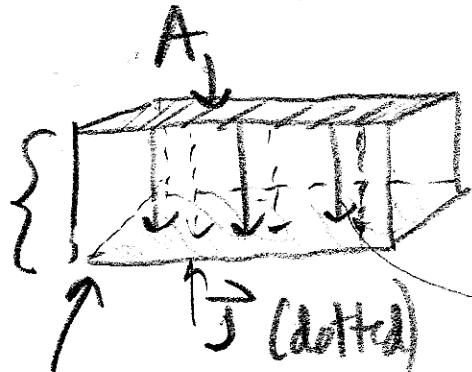
a constant known as conductivity; resistivity

the larger σ , the more current density for a given \vec{E} applied.

$$\vec{J} = \sigma \vec{E} \quad \text{"intrinsic quantities"}$$

leads to

$$V = IR. \quad \text{"extrinsic"}$$



σ in material

$$E = \frac{V}{l}, \quad J = \sigma E = \frac{\sigma V}{l}$$

$$I = J \cdot A = \frac{\sigma V A}{L}$$

$$V = \left(\frac{L}{\sigma A}\right) \cdot I$$

$$R = \boxed{\frac{L}{\sigma A}}$$

P (new quantity)
not charge
density

$$= \frac{1}{\sigma}$$

$$\boxed{R = \frac{L}{A} \cdot P = \frac{L}{A} \frac{1}{\sigma}}$$

units:

CGS

$$\vec{J} = \sigma \vec{E}$$

\uparrow \uparrow \nwarrow
 $\frac{esu}{cm^2 \cdot s}$ $\frac{1}{s}$ $\frac{esu}{cm^2}$
 $(! ! ! !)$

SI/MKS

$$R = \frac{L}{A} \cdot P$$

\uparrow \uparrow \uparrow
 ohms (SI) = $\frac{m}{m^2} \cdot (\text{ohm} \cdot m)$
 or $(\text{ohm} \cdot \text{cm})$

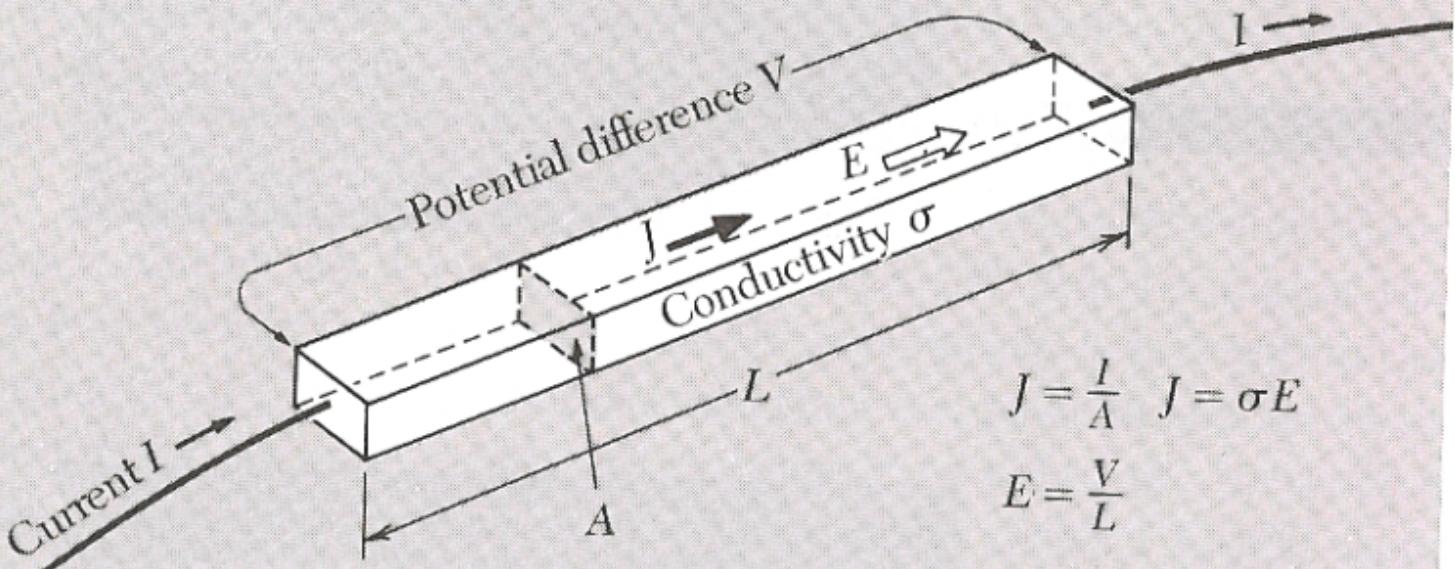
There are some conditions for ohm's law... see text

TABLE 4.1

Resistivity and its reciprocal, conductivity, for a few materials

Material	Resistivity ρ	Conductivity σ
Pure copper, 273 K	1.56×10^{-6} ohm-cm 1.73×10^{-18} sec	6.4×10^5 (ohm-cm) $^{-1}$ 5.8×10^{17} sec $^{-1}$
Pure copper, 373 K	2.24×10^{-6} ohm-cm 2.47×10^{-18} sec	4.5×10^5 (ohm-cm) $^{-1}$ 4.0×10^{17} sec $^{-1}$
Pure germanium, 273 K	200 ohm-cm 2.2×10^{-10} sec	0.005 (ohm-cm) $^{-1}$ 4.5×10^9 sec $^{-1}$
Pure germanium, 500 K	0.12 ohm-cm 1.3×10^{-13} sec	8.3 (ohm-cm) $^{-1}$ 7.7×10^{12} sec $^{-1}$
Pure water, 291 K	2.5×10^7 ohm-cm 2.8×10^{-5} sec	4.0×10^{-8} (ohm-cm) $^{-1}$ 3.6×10^4 sec $^{-1}$
Seawater (varies with salinity)	25 ohm-cm 2.8×10^{-11} sec	0.04 (ohm-cm) $^{-1}$ 3.6×10^{10} sec $^{-1}$

Note: 1 ohm-meter = 100 ohm-cm = 1.11×10^{-10} sec.

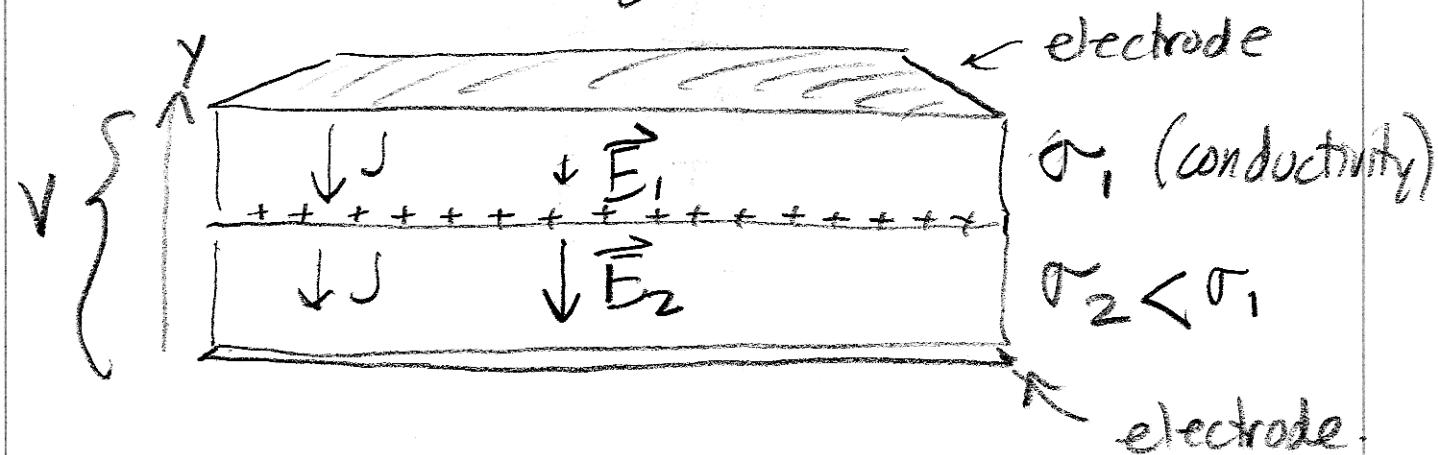


$$J = \frac{I}{A} \quad J = \sigma E$$

$$E = \frac{V}{L}$$

$$R = \frac{V}{I} = \frac{L}{\sigma A}$$

One interesting example:



What is the same? In, when steady state.

$$\nabla \cdot \vec{J} = \frac{\partial J_x}{\partial x} + \frac{\partial J_y}{\partial y} + \frac{\partial J_z}{\partial z} = -\frac{\partial \phi}{\partial x}$$

(∞ in those directions)

In steady state.

$$\frac{\partial J_x}{\partial y} = 0 \rightarrow J_y \text{ is constant}$$

What is different?

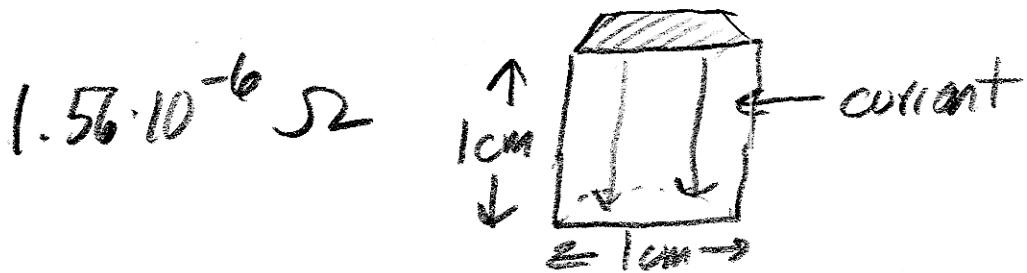
$$\vec{E}_1 = \left(\frac{1}{\sigma_1}\right) \vec{J} < \left(\frac{1}{\sigma_2}\right) \vec{J} < \vec{E}_2$$

(new field lines must start)

Must be some positive charge at the interface... in steady state, Initially not (then, $\frac{\partial \phi}{\partial t} \neq 0$), but it builds up.

Interesting Table - p. 133

$\rho = \frac{1}{\sigma}$: Pure Copper .. $1.56 \cdot 10^{-6}$ ohm-cm



Pure Water : $25 \cdot 10^6$ ohm-cm

$25 M\Omega$ above.

Sea Water : 25 ohm-cm
(Bath Water)

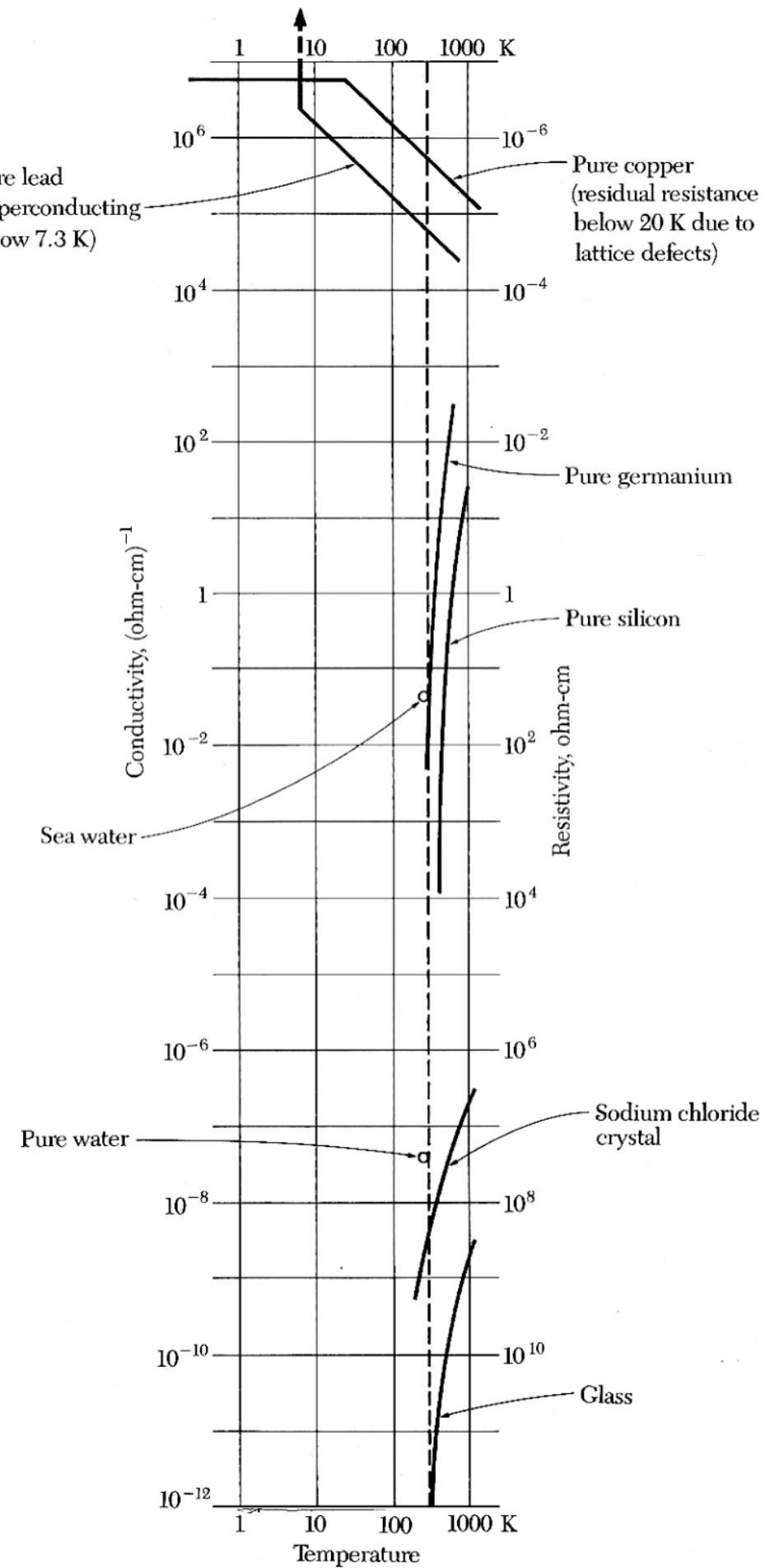
purity matters a lot!

Conductivity/Resistivity are very sensitive to temperature.. look at plot on page 140.

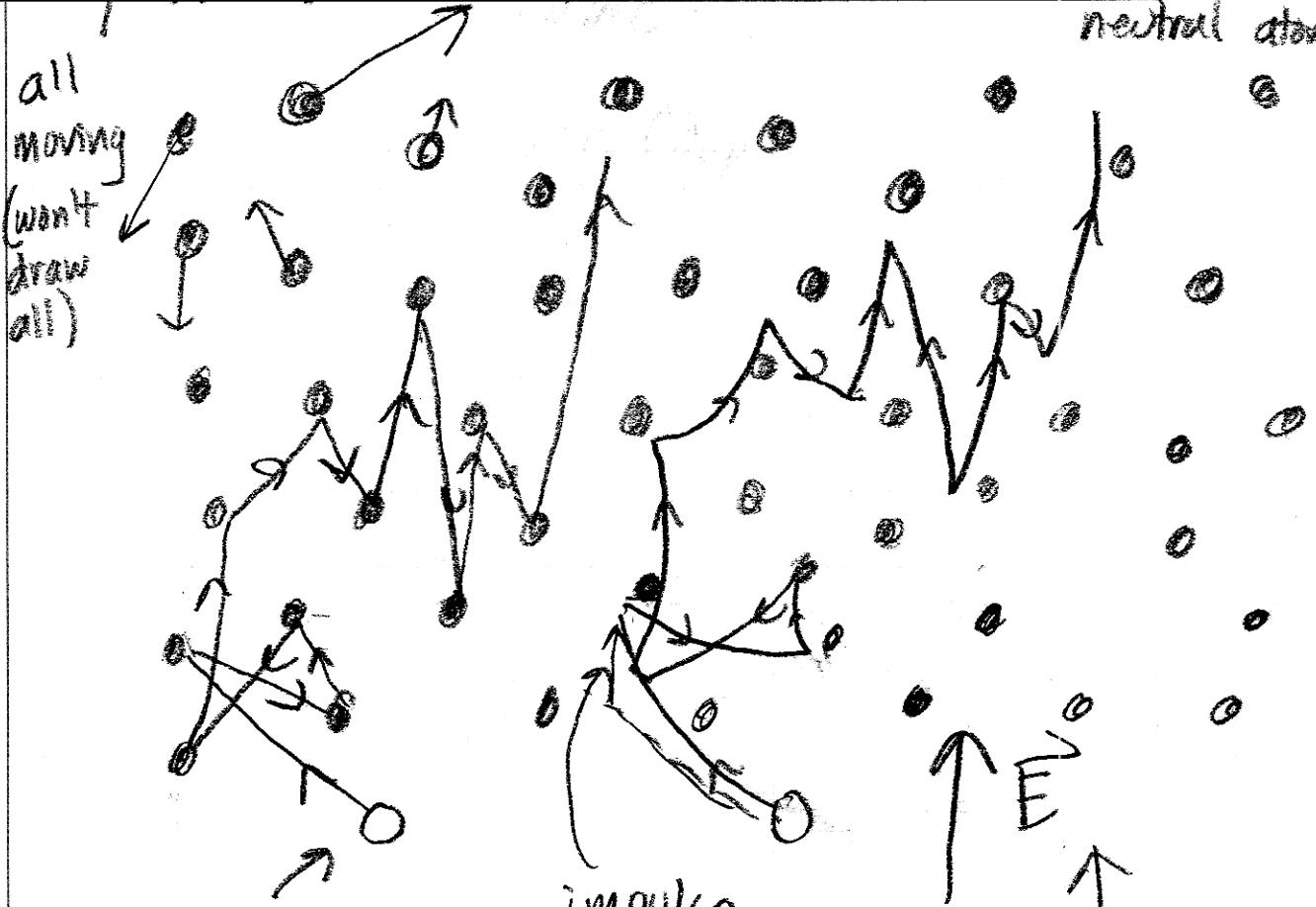
Reason: dominant motion¹ in condensed matter is thermal motion, even when electric field is present.

Electric field is a minor effect.

Usually.



neutral atoms



ion, q ,
 n_q , M_q
 NO ELECTRIC
 FIELD

impulse
 due to
 field.

bunch of neutral
 atoms in
 solid, liquid

$$\langle \vec{v} \rangle = 0$$

$$\frac{1}{2} m \langle v^2 \rangle = \frac{3}{2} kT$$

$$\langle v^2 \rangle = \frac{3kT}{m}$$

$$\langle v^2 \rangle^{1/2} = \sqrt{\frac{3kT}{M}}$$

suppose: $T \approx 300 \text{ K}$, $m \approx 0.030 \text{ kg}$

$$\langle \Delta v \rangle = \frac{qE}{M} \langle \Delta t \rangle$$

$$\langle \vec{v} \rangle = \langle \vec{\Delta v} \rangle$$

$$\langle v^2 \rangle^{1/2} = \sqrt{\frac{3 \cdot 8.3 \cdot 300}{0.030}} = 500 \frac{\text{m}}{\text{s}}$$

big!