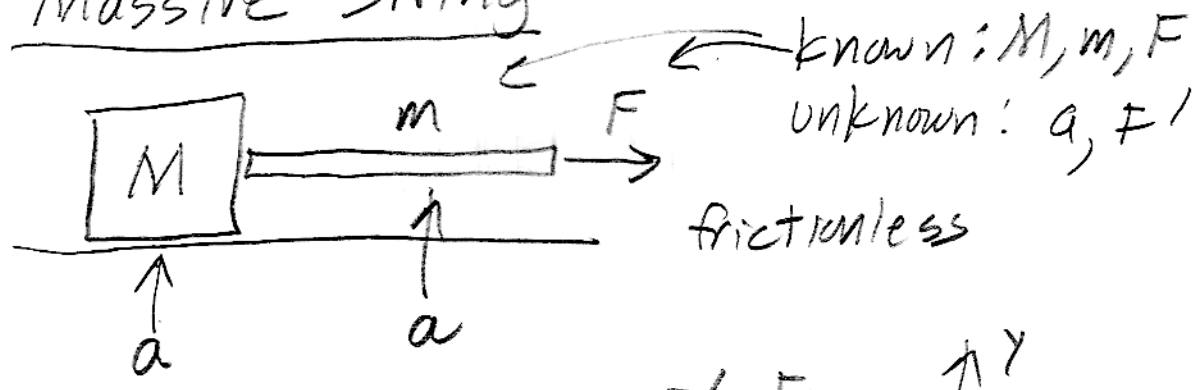
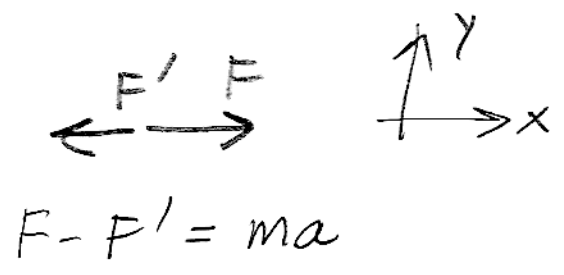


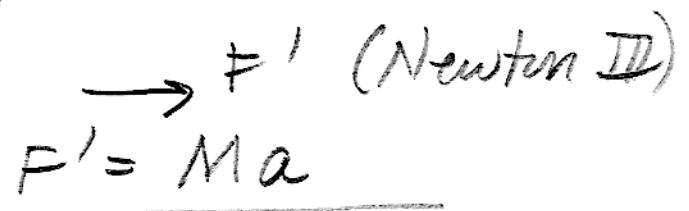
Massive String



① the string



② the mass M



$$F = (m + M)a$$

$$a = \frac{F}{m + M}$$

$m \ll M$
 $a \approx F/M$

$$F - F' = \frac{F}{m + M} \times m$$

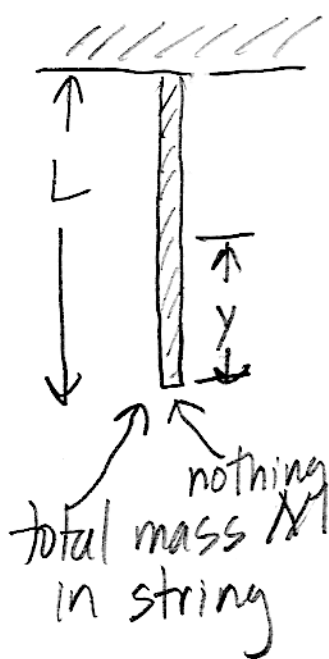
$$F \left(1 - \frac{m}{m + M}\right) = F'$$

$$F' = \frac{M}{m + M} F$$

$m \ll M$
 $F' \approx F$

If string is uniform in mass density.





assuming mass density uniform

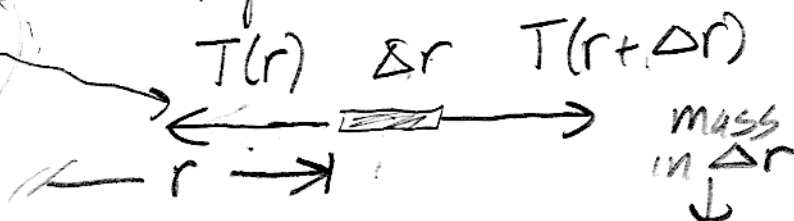
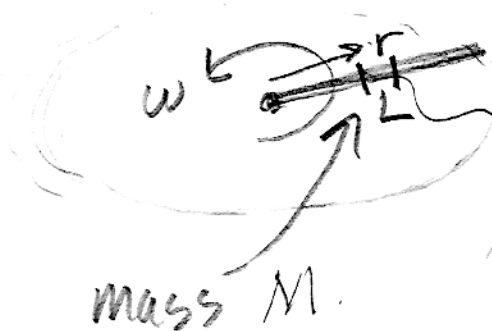
mass in an interval y
 $= \left(\frac{y}{L}\right) M$

Draw forces on that mass in the interval:

$\uparrow T(y)$
 $T(y) - \left(\frac{y}{L}\right) Mg = 0$
 $\downarrow \left(\frac{y}{L}\right) Mg$
 $T(y) = \left(\frac{y}{L}\right) Mg$

Whirling Rope (big fun)

Small segment of rope:



$T(r+\Delta r) - T(r) = \left(\frac{M}{L} \Delta r\right) \cdot (-a)$
 centripetal acceleration
 $= \omega^2 r$

$\Delta T = \frac{-M r \omega^2}{L} \cdot \Delta r$

calculus:
 $\frac{dT}{dr} = \frac{-M r \omega^2}{L}$

Ignore gravity
mass density: $\frac{M}{L}$

$$T(r) = -\frac{1}{2} \frac{M\omega^2}{L} r^2 + \text{constant}$$

constant: when $r=L$, $T(L)=0$
nothing pulling on the last bit of rope

$$\text{so, } T(L)=0 = -\frac{1}{2} \frac{M\omega^2}{L} \cdot L^2 + \text{constant}$$

$$\text{constant} = \frac{1}{2} \frac{M\omega^2}{L} L^2$$

$$T(r) = \frac{1}{2} \frac{M\omega^2}{L} (L^2 - r^2)$$

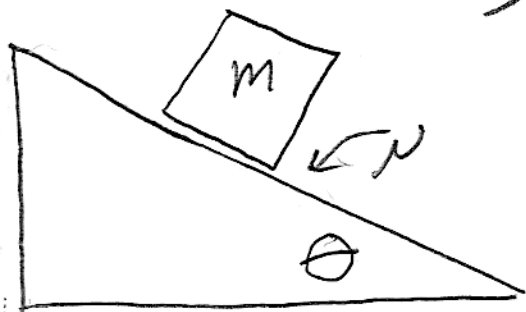
Friction

- depends on Normal force
- \perp to Normal force.
- resists applied force.
- when body at rest,

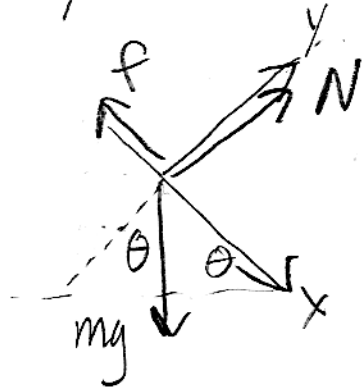
$$f = \text{friction force} \quad 0 \leq f \leq \mu N$$

and f is always unknown.

- when body no longer at rest
 $f = \mu N$, opposes motion.



as θ increases, when does the mass slip?



$$y: N - mg \cos \theta = 0$$

$$N = mg \cos \theta$$

note: as $\theta \uparrow$ $N \downarrow$

$$x: -f + mg \sin \theta = 0 \text{ (max)}$$

↑ sticking ↑ sliding

which?

assume $f = \mu N$

remember: friction only "brakes" never pushes

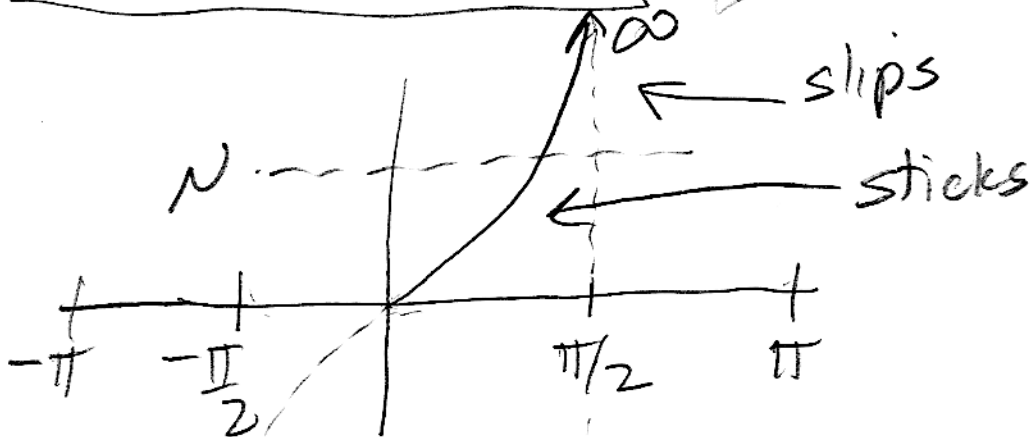
$$-\mu N + mg \sin \theta \geq 0$$

$$-\mu mg \cos \theta + mg \sin \theta \geq 0$$

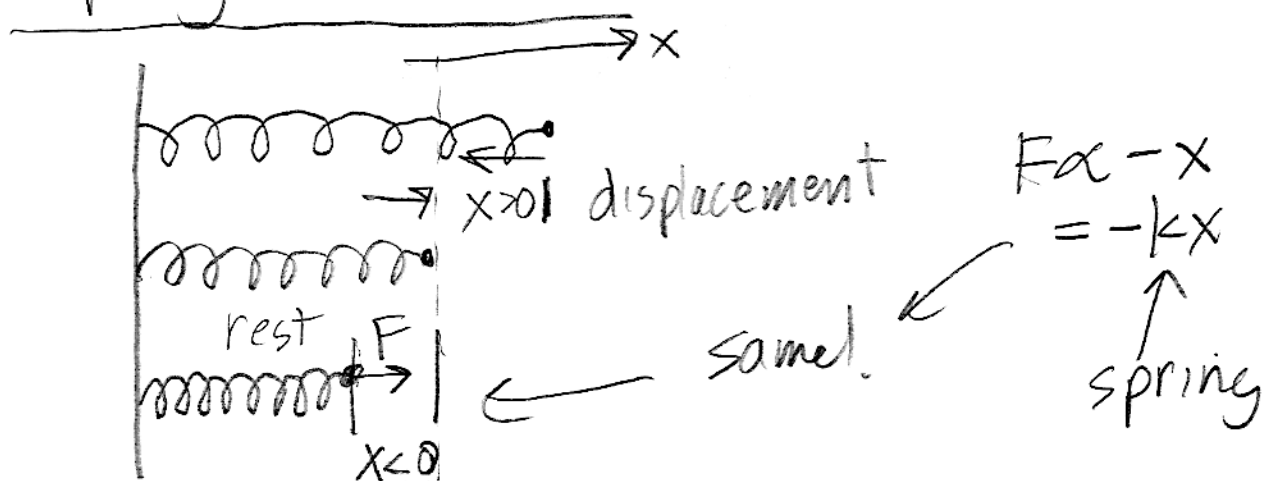
$$\sin \theta \geq \mu \cos \theta$$

$\tan \theta \geq \mu$
for slipping

$\tan \theta < \mu$ sticking



Springs - first look



now $-kx = ma_x = m\ddot{x} = m\frac{d^2x}{dt^2}$

or $\ddot{x} + \left(\frac{k}{m}\right)x = 0$

try $x = A\sin(\omega t) + B\cos(\omega t)$

$$\dot{x} = +\omega A\cos(\omega t) - \omega B\sin(\omega t)$$

$$\ddot{x} = -\omega^2 A\sin(\omega t) - \omega^2 B\cos(\omega t)$$

$$= -\omega^2 (A\sin(\omega t) + B\cos(\omega t))$$

$$\ddot{x} = -\omega^2 x$$

$$\ddot{x} + \omega^2 x = 0$$

works as long as $\omega^2 = \frac{k}{m}$

$$\omega = \sqrt{\frac{k}{m}}$$