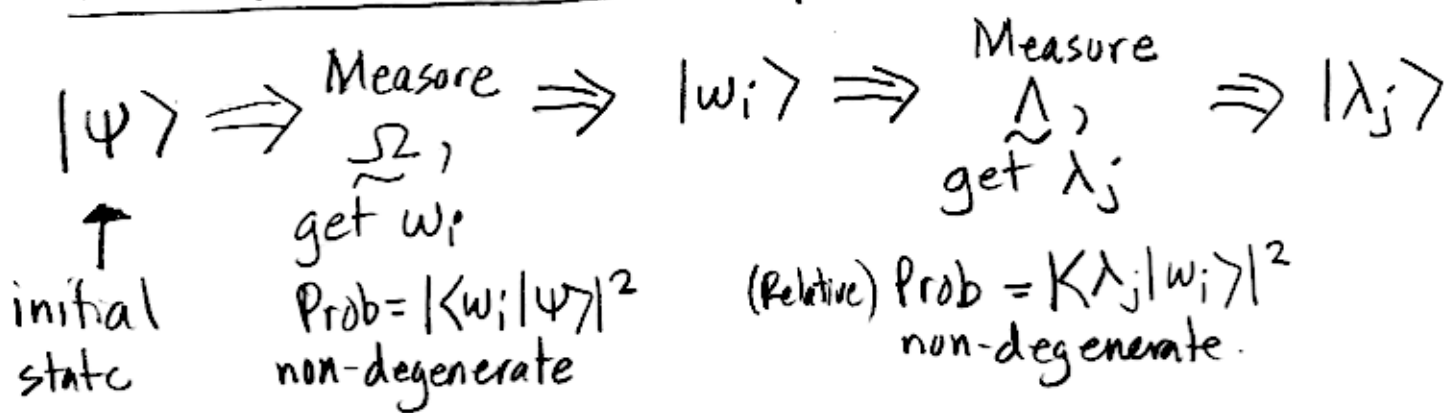


Compatible + Incompatible Variables p.129 80



Discussion : Suppose we repeat measurements many times, always starting from the same (or identically prepared) initial state $|\psi\rangle$

Points :

(1) Can we always get same ^{particular} w_i ?
In general, no, will get a variety of w_i 's, each with probability $P(w_i) = |\langle w_i | \psi \rangle|^2$. Of course, when $|\psi\rangle$ itself is an eigenstate, then we'll always get the corresponding eigenvalue.

(2) Given that a particular w_i results from the first measurement, when will a specific λ_j always result from the second measurement?

\Rightarrow when $|\langle \lambda_j | w_i \rangle|^2 = 1$

or $|w_i\rangle$ also eigenket of $\hat{\Lambda}_j$
with eigenvalue λ_j

$$\text{so } \hat{\Omega} |w; \lambda_j\rangle = w_i |w; \lambda_j\rangle$$

$$\text{and } \hat{\Lambda} |w; \lambda_j\rangle = \lambda_j |w; \lambda_j\rangle$$

$$\text{further } \hat{\Lambda} \hat{\Omega} |w; \lambda_j\rangle = \hat{\Lambda} w_i |w; \lambda_j\rangle = w_i \hat{\Lambda} |w; \lambda_j\rangle = w_i \lambda_j |w; \lambda_j\rangle$$

$$\text{and } \hat{\Omega} \hat{\Lambda} |w; \lambda_j\rangle = \hat{\Omega} \lambda_j |w; \lambda_j\rangle = \lambda_j \hat{\Omega} |w; \lambda_j\rangle = \lambda_j w_i |w; \lambda_j\rangle$$

$$\text{so } (\hat{\Omega} \hat{\Lambda} - \hat{\Lambda} \hat{\Omega}) |w; \lambda_j\rangle = 0 \quad \star$$

How can we get this?

(A) Obvious: $[\hat{\Omega}, \hat{\Lambda}] = 0$ always

then... $\hat{\Omega} + \hat{\Lambda}$ are compatible.

Converse....

(B) $[\hat{\Omega}, \hat{\Lambda}] =$ something obviously non-zero
"Incompatible"

example: \hat{x} and $\hat{p} = \hbar \hat{k}$, $\hat{k} \equiv \frac{1}{i} \frac{d}{dx}$

$$\text{so, } [\hat{x}, \hat{p}] |f\rangle \equiv [x, \frac{\hbar}{i} \frac{d}{dx}] f(x)$$

$$\equiv \left(x \frac{\hbar}{i} \frac{d}{dx} - \frac{\hbar}{i} \frac{d}{dx} x \right) f(x)$$

$$\equiv x \frac{\hbar}{i} \frac{df}{dx} - \frac{\hbar}{i} \left(f(x) + x \frac{df}{dx} \right)$$

$$[\hat{x}, \hat{p}] |f\rangle \equiv -\frac{\hbar}{i} f(x)$$

$$\text{or } \boxed{[\hat{x}, \hat{p}] = -\frac{\hbar}{i} = i\hbar}$$

famous
incompatible
variables

... origin of the Heisenberg uncertainty principle...

Example

$$\underline{\Omega} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad \underline{\Lambda} = \begin{pmatrix} a & b & d \\ b^* & c & b^* \\ d & b & a \end{pmatrix}$$

(Unitary...)
aka, "symmetry operator"

$$\underline{\Omega} \underline{\Lambda} - \underline{\Lambda} \underline{\Omega} \stackrel{?}{=} 0$$

$$\begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} a & b & d \\ b^* & c & b^* \\ d & b & a \end{pmatrix} = \begin{pmatrix} d & b & a \\ b^* & c & b^* \\ a & b & d \end{pmatrix}$$

$$\begin{pmatrix} a & b & d \\ b^* & c & b^* \\ d & b & a \end{pmatrix} \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} = \begin{pmatrix} d & b & a \\ b^* & c & b^* \\ a & b & d \end{pmatrix}$$

but note: $\underline{\Omega} \underline{\Lambda} - \underline{\Lambda} \underline{\Omega} = 0$

$$\underline{\Lambda}^{-1} \underline{\Omega} \underline{\Lambda} - \underline{\Omega} = 0$$

$$\underline{\Lambda}^+ \underline{\Omega} \underline{\Lambda} = \underline{\Omega}$$

$$\underline{\Lambda}^{-1} = \underline{\Lambda}^+ \text{ (Unitary)}$$

← $\underline{\Omega}$ is "invariant" under $\underline{\Lambda}$.

$\underline{\Lambda}$ is a "symmetry operator"

$$\begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = 1 \cdot \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

obvious eigenvector.

$$\begin{vmatrix} \lambda & 0 & 1 \\ 0 & 1-\lambda & 0 \\ 1 & 0 & \lambda \end{vmatrix} = \lambda \begin{vmatrix} 1-\lambda & 0 \\ 0 & \lambda \end{vmatrix} + 1 \begin{vmatrix} 0 & 1-\lambda \\ 1 & 0 \end{vmatrix} = 0$$

$$\lambda(1-\lambda)\lambda - (1-\lambda) = 0$$

$$(\lambda^2 - 1)(1-\lambda) = 0 \quad \lambda = \pm 1, 1$$

given $|w_3\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$ $w_3 = 1$

$|w_1\rangle, |w_2\rangle$ of form $\begin{pmatrix} \alpha \\ 0 \\ \beta \end{pmatrix}$

$$\pm 1: \begin{pmatrix} \mp 1 & 0 & 1 \\ 0 & 1 \mp 1 & 0 \\ 1 & 0 & \mp 1 \end{pmatrix} \begin{pmatrix} \alpha \\ 0 \\ \beta \end{pmatrix} \quad \begin{matrix} +1: \alpha + \beta = 0, \\ -1: \alpha + \beta = 0, \end{matrix} \quad \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix}$$

$$\begin{pmatrix} 1/\sqrt{2} \\ 0 \\ -1/\sqrt{2} \end{pmatrix}$$

A transformation that diagonalizes $\underline{\Omega}$ is:

$$\begin{pmatrix} 1/\sqrt{2} & 0 & -1/\sqrt{2} \\ 1/\sqrt{2} & 0 & 1/\sqrt{2} \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} & 1/\sqrt{2} & 0 \\ 0 & 0 & 1 \\ -1/\sqrt{2} & 1/\sqrt{2} & 0 \end{pmatrix}$$

$$\begin{matrix} \underline{U}^\dagger & \underline{\Omega} \\ \begin{pmatrix} 1/\sqrt{2} & 0 & -1/\sqrt{2} \\ 1/\sqrt{2} & 0 & 1/\sqrt{2} \\ 0 & 1 & 0 \end{pmatrix} & \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{matrix} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

now!

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$$|\psi\rangle \Rightarrow \text{Measure } \hat{\Omega} \Rightarrow \begin{matrix} -1 : |\omega_1\rangle \\ +1 : \frac{(\langle\omega_2|\psi\rangle + \langle\omega_3|\psi\rangle)}{\sqrt{|\langle\omega_2|\psi\rangle|^2 + |\langle\omega_3|\psi\rangle|^2}} \end{matrix} \Rightarrow \text{Measure } \hat{\Lambda}$$

$$-1, \text{ prob } |\langle\omega_1|\psi\rangle|^2$$

$$= \left| \begin{pmatrix} 1/\sqrt{2} & 0 & -1/\sqrt{2} \\ 1/\sqrt{2} & 0 & 1/\sqrt{2} \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \right|^2$$

$$+1, \text{ prob: a) } 1 - |\langle\omega_1|\psi\rangle|^2$$

$$b) |\langle\omega_2|\psi\rangle|^2 + |\langle\omega_3|\psi\rangle|^2$$

+1 is ambiguous!

Transform $\hat{\Lambda}$ into eigenbasis of $\hat{\Omega}$:

$$\begin{pmatrix} 1/\sqrt{2} & 0 & -1/\sqrt{2} \\ 1/\sqrt{2} & 0 & 1/\sqrt{2} \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} a & b & d \\ b^* & c & b^* \\ d & b & a \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} & 1/\sqrt{2} & 0 \\ 0 & 0 & 1 \\ -1/\sqrt{2} & 1/\sqrt{2} & 0 \end{pmatrix}$$

$$\begin{pmatrix} 1/\sqrt{2} & 0 & -1/\sqrt{2} \\ 1/\sqrt{2} & 0 & 1/\sqrt{2} \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}}(a-d) & \frac{1}{\sqrt{2}}(a+d) & b \\ 0 & \sqrt{2} b^* & c \\ -\frac{1}{\sqrt{2}}(a-d) & \frac{1}{\sqrt{2}}(a+d) & b \end{pmatrix}$$

$$\hat{\Lambda} \equiv \begin{pmatrix} a-d & 0 & 0 \\ 0 & a+d & \sqrt{2} b \\ 0 & \sqrt{2} b^* & c \end{pmatrix}$$

← measuring $\hat{\Lambda}$ can "break" degeneracy...

- need only diagonalize 2×2 piece...
- any linear combinations of $|w_2\rangle$ and $|w_3\rangle$ are still eigenkets of $\hat{\Omega}$ with eigenvalue $w_2 = w_3 = +1$

$$\begin{vmatrix} a+d-\lambda & \sqrt{2}b^* \\ \sqrt{2}b^* & c-\lambda \end{vmatrix} = (a+d-\lambda)(c-\lambda) - 2|b|^2 = 0$$

$$\lambda^2 - (a+c+d)\lambda + c(a+d) - 2|b|^2 = 0$$

Discriminant: $(a+c+d)^2 - 4c(a+d) + 8|b|^2$

$$= (a+d)^2 + (c)^2 + 2c(a+d) - 4c(a+d) + 8|b|^2$$

$$= (a+d-c)^2 + 8|b|^2$$

so $\lambda = \frac{(a+c+d) \pm [(a+d-c)^2 + 8|b|^2]^{1/2}}{2}$

$$\lambda = \frac{1}{2}(a+d+c) \pm \left[\left(\frac{a+d-c}{2} \right)^2 + (\sqrt{2}|b|)^2 \right]^{1/2}$$

Midterm: $a = \frac{2}{3}b$ $c = b$ $d = \frac{1}{3}b$ $b^* = b$

$$\lambda = \frac{1}{2} \left(\frac{2}{3} + \frac{1}{3} + 1 \right) b \pm \left[\left(\frac{\frac{2}{3} + \frac{1}{3} - 1}{2} \right)^2 b^2 + (\sqrt{2}b)^2 \right]^{1/2}$$

$$\lambda = b \pm \sqrt{2}b$$

Eigenvectors: $\lambda = b + \sqrt{2}b$: $\begin{pmatrix} \frac{1}{3}b - b - \sqrt{2}b & 0 & 0 \\ 0 & -\sqrt{2}b & \sqrt{2}b \\ 0 & \sqrt{2}b & -\sqrt{2}b \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$

$$\alpha = \beta,$$

$$\begin{pmatrix} 0 \\ 1/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix}$$

representation
in $\underline{\Omega}$ eigenbasis

meaning $\frac{1}{\sqrt{2}} \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix} + \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1/2 \\ 1/\sqrt{2} \\ 1/2 \end{pmatrix}$ in original basis

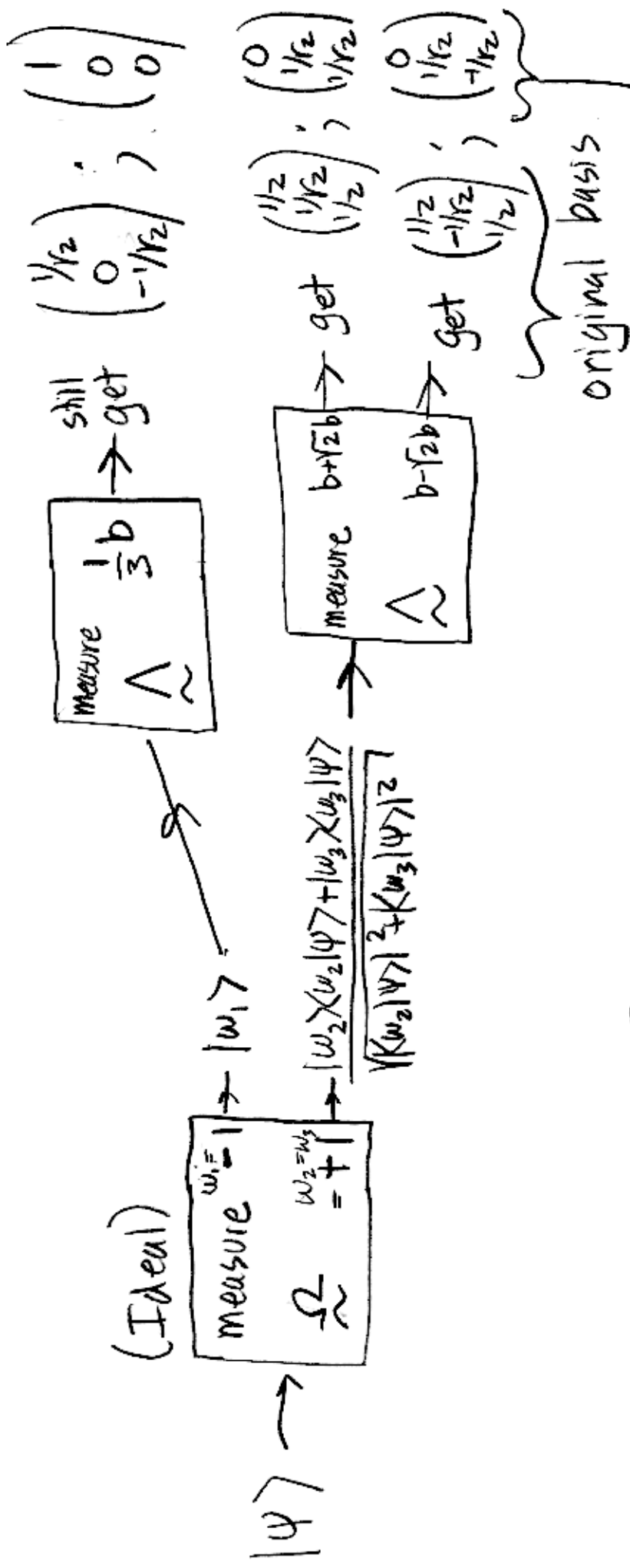
$$\lambda = b + \sqrt{2}b$$

$$\lambda = b - \sqrt{2}b : \begin{pmatrix} \frac{1}{3}b - b + \sqrt{2}b & 0 & 0 \\ 0 & \sqrt{2}b & \sqrt{2}b \\ 0 & \sqrt{2}b & \sqrt{2}b \end{pmatrix} \begin{pmatrix} 0 \\ \alpha \\ \beta \end{pmatrix} = 0$$

$$\alpha = -\beta, \quad \begin{pmatrix} 0 \\ 1/\sqrt{2} \\ -1/\sqrt{2} \end{pmatrix}$$

representation
in $\underline{\Omega}$ eigenbasis.

meaning $\frac{1}{\sqrt{2}} \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix} - \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1/2 \\ -1/\sqrt{2} \\ 1/2 \end{pmatrix}$ in original basis



Since there are 3 distinct combinations of eigenvalues:

$$(-1, \frac{1}{3}b)$$

$$(+1, b + \sqrt{2}b)$$

$$(+1, b - \sqrt{2}b)$$

Ω_2, Λ form a "Complete Set of Commuting Observables"

still eigenkets of Ω_2 , eigenvalue +1

first eigenbasis of Ω_2